

**Domino-tilings of Aztec diamonds,
Random Surfaces
& the Gaussian Unitary Ensemble (GUE)**

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DYGEST Conference at Corfou (May 2013)

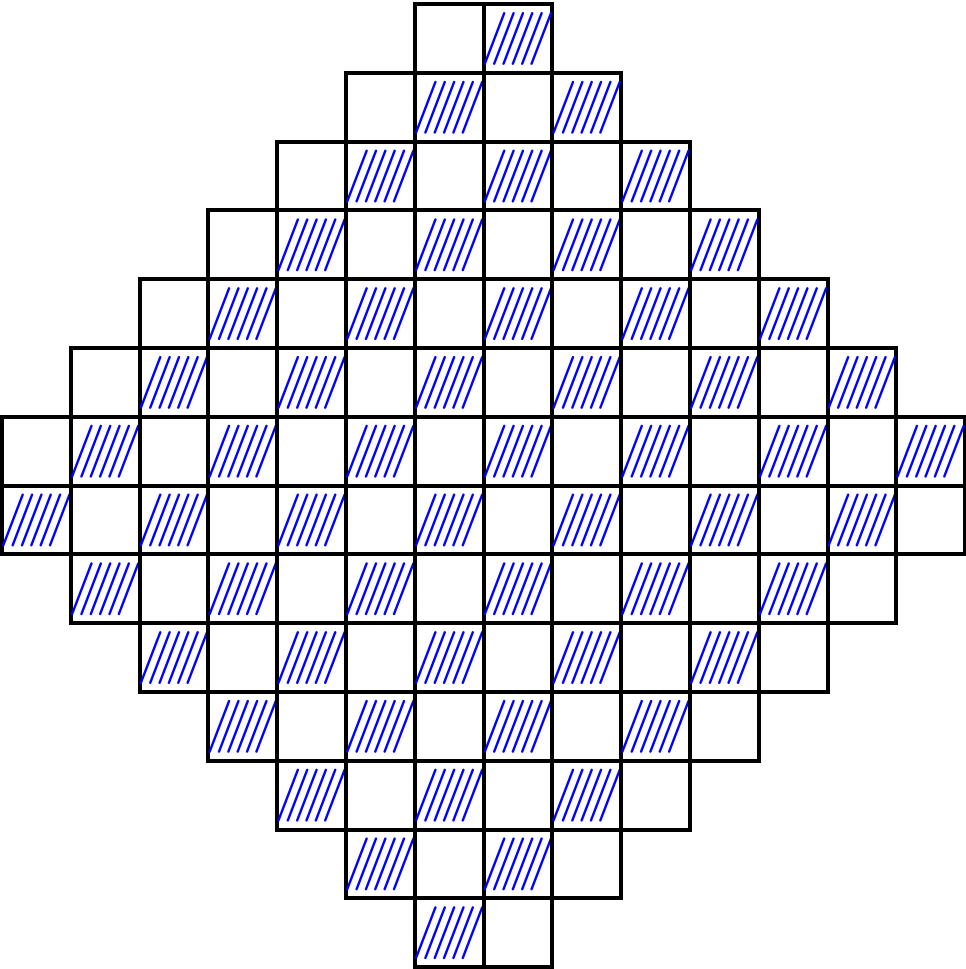
1. Domino tilings of Aztec Diamonds & Random Surfaces

2. Domino tilings of double Aztec Diamonds

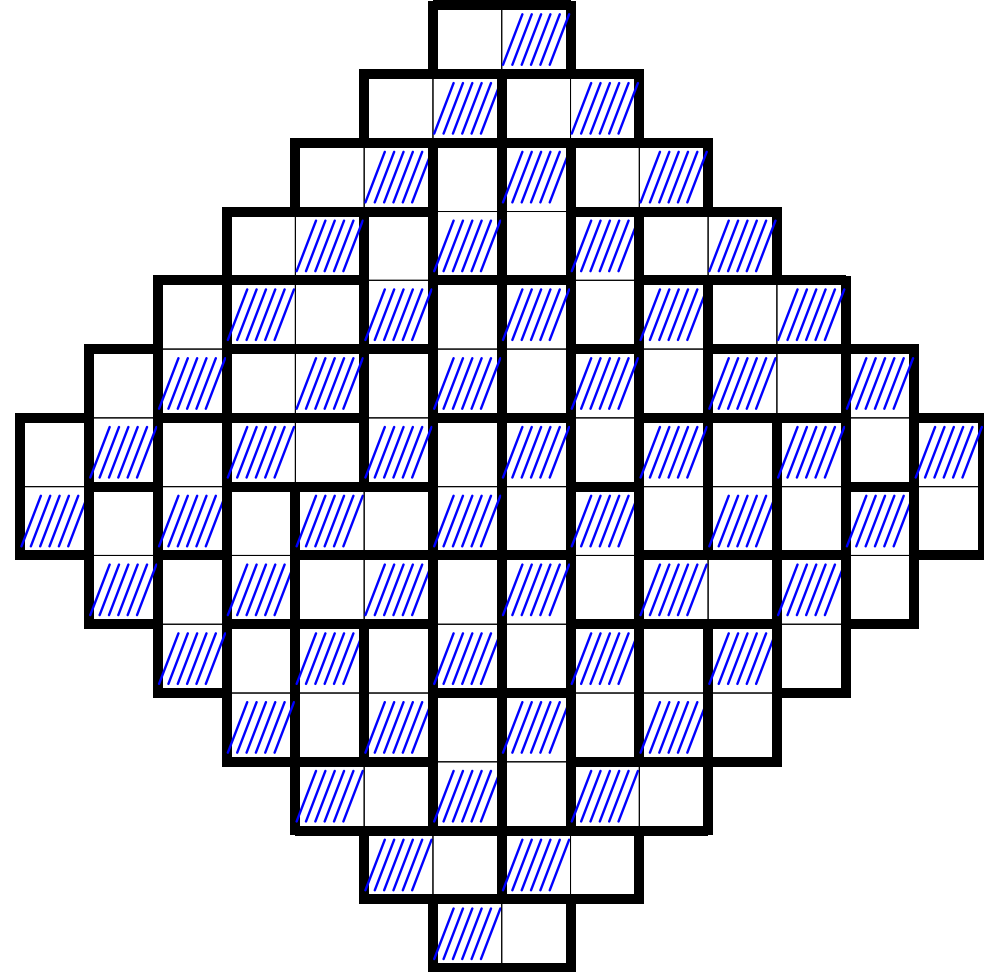
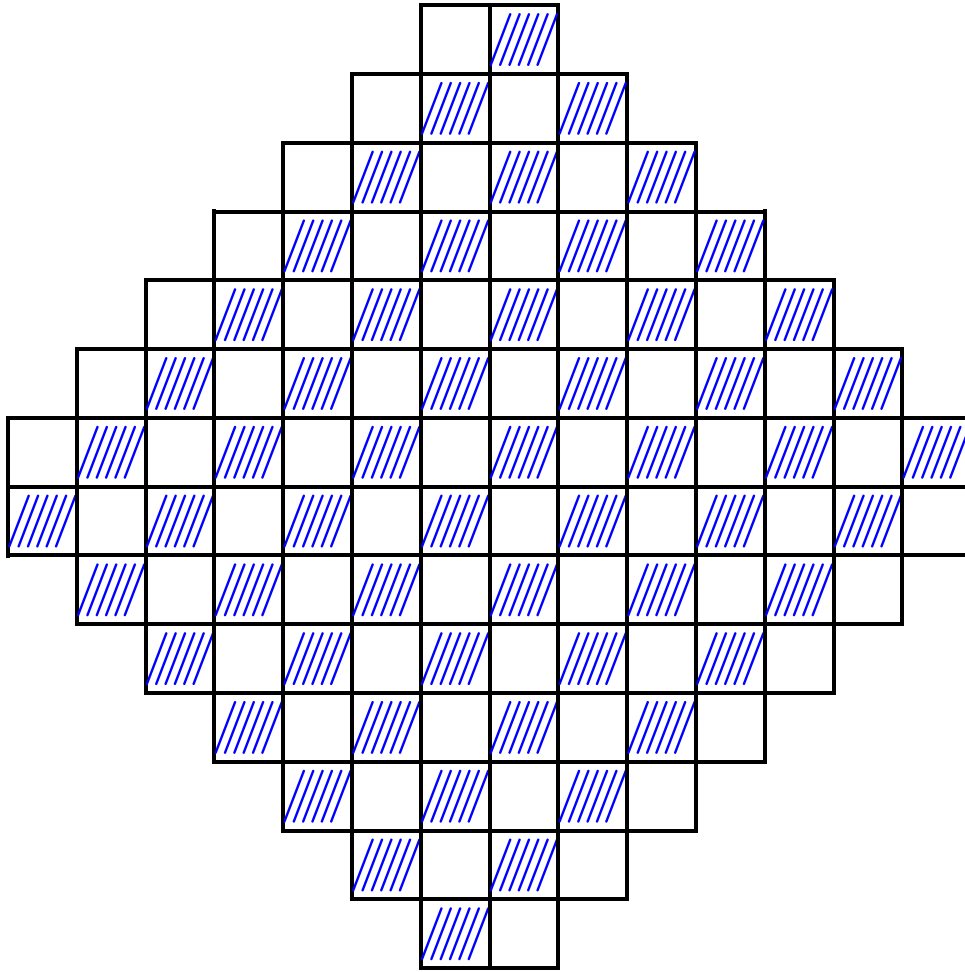
3. Letting the size $n \rightarrow \infty$

1. Domino tilings of Aztec Diamonds & Random Surfaces

An Aztec diamond (size n):



Domino tilings of an Aztec diamond :

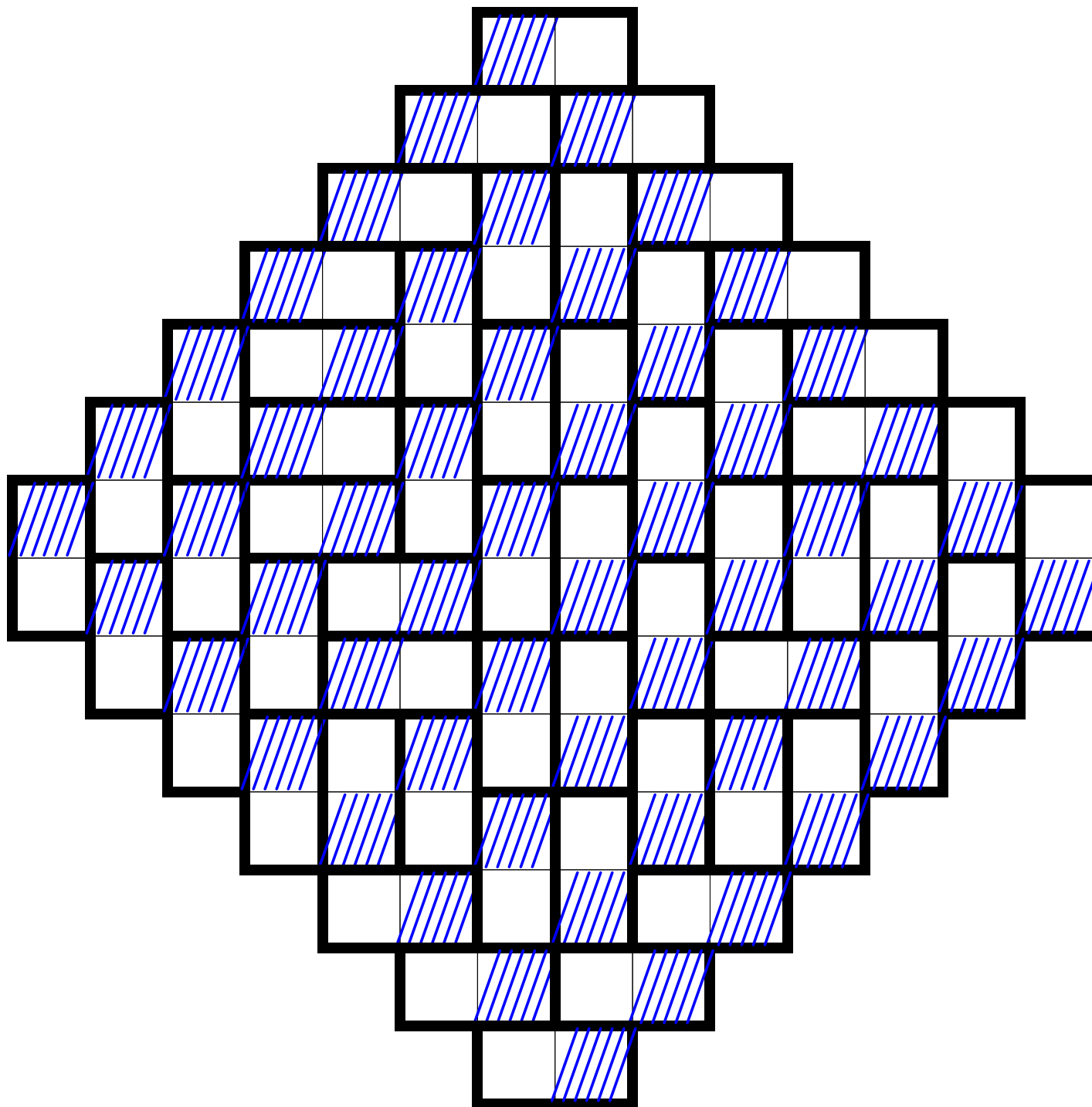


4 types of situations for the domino's!

Kasteleyn 1961, Elkies, Kuperberg, Larsen, Propp '92

Cohn, Elkies, Propp '96, Johansson '00, '03, '05

M. Fulmek, C. Krattenthaler '00, Johansson-Nordenstam '06



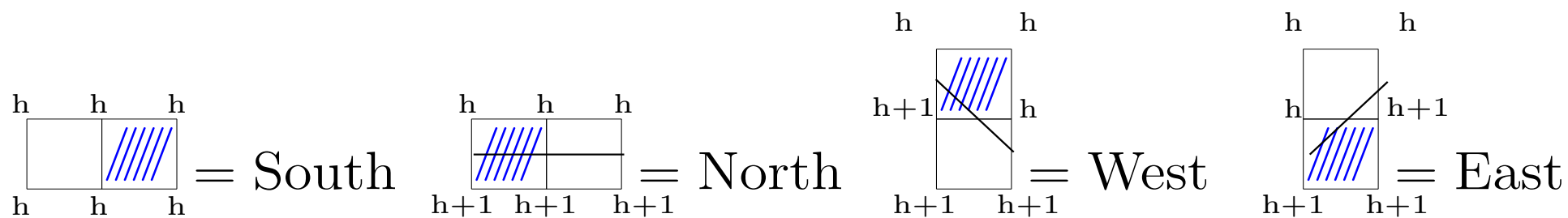
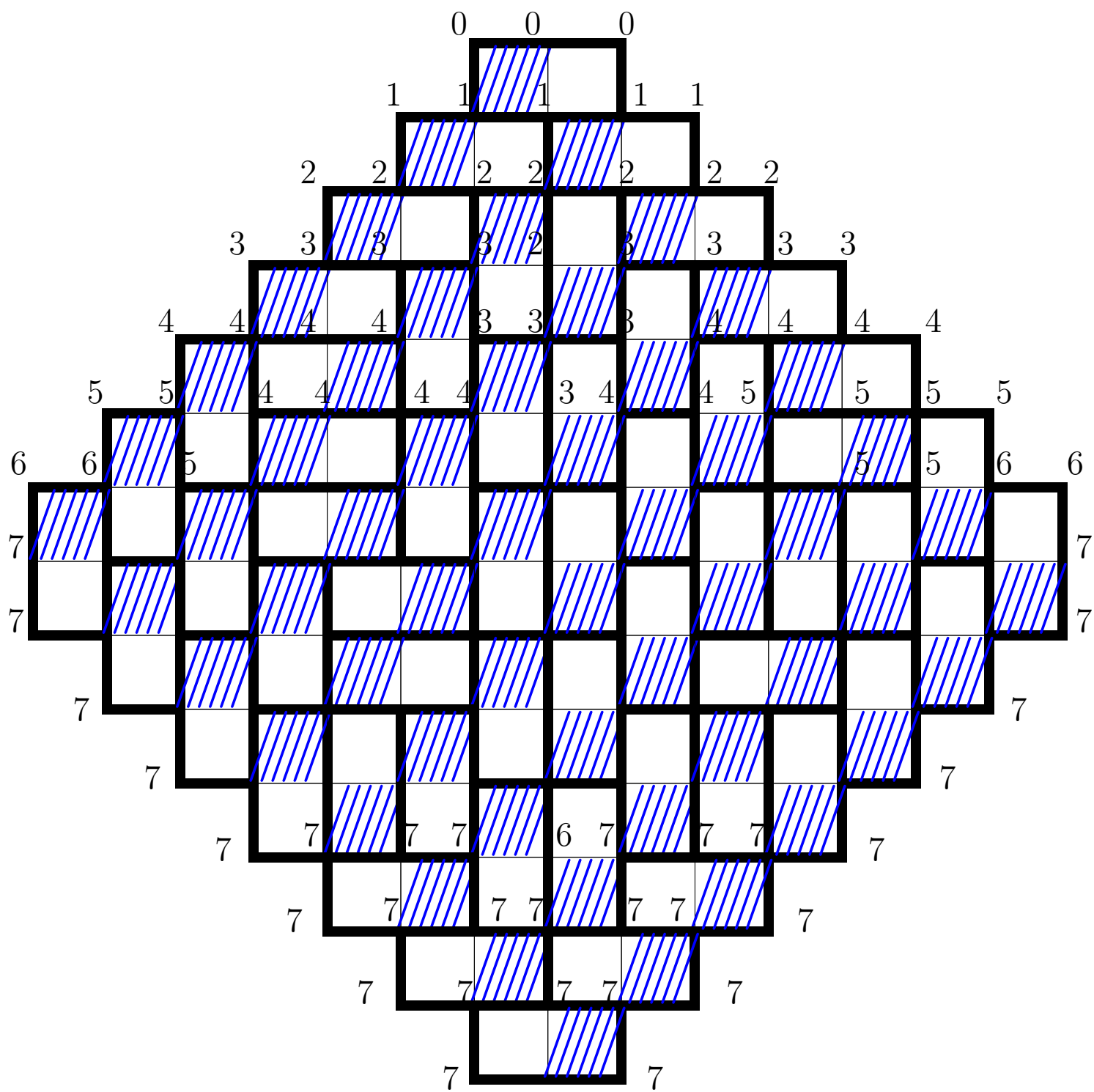
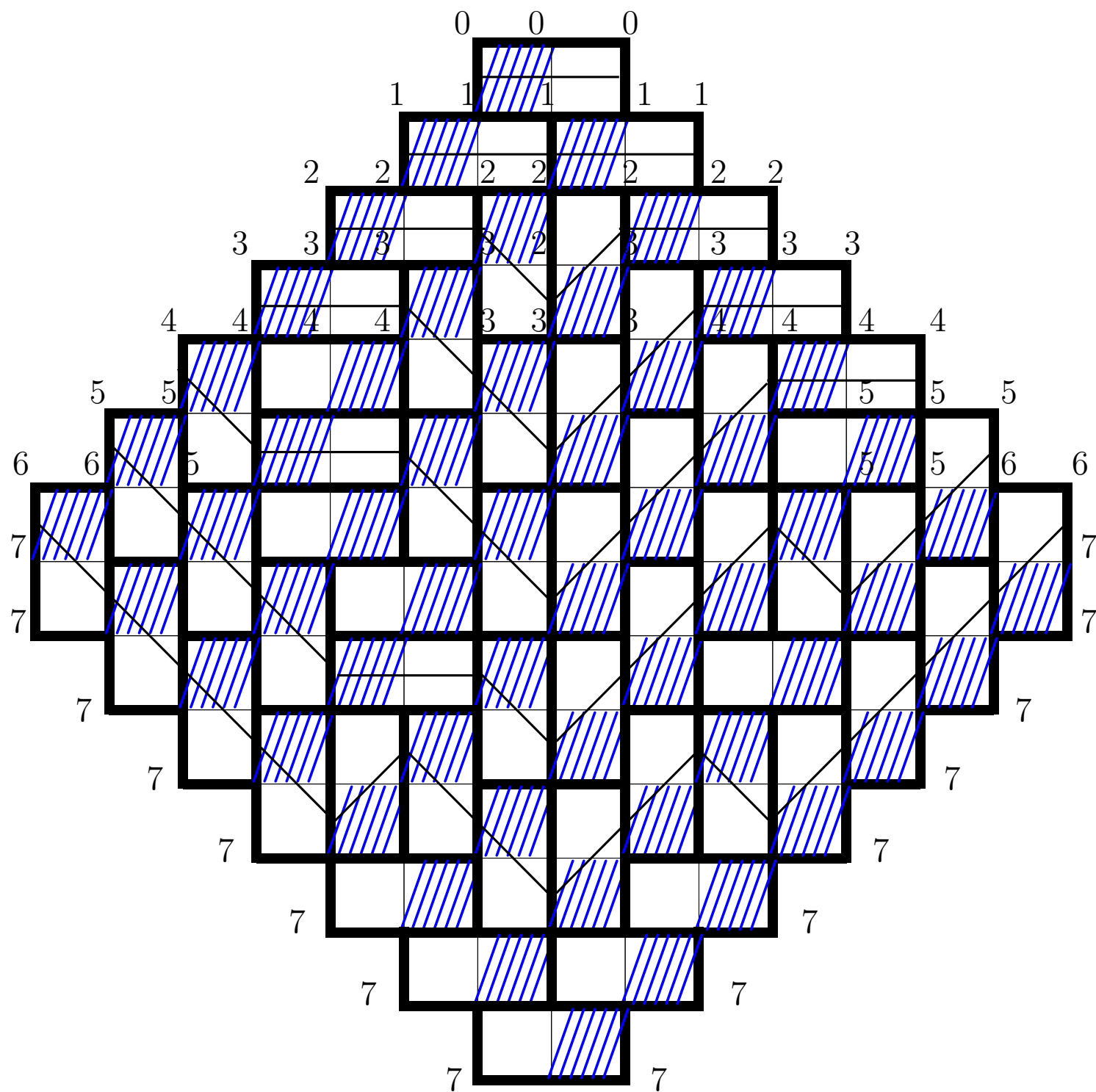
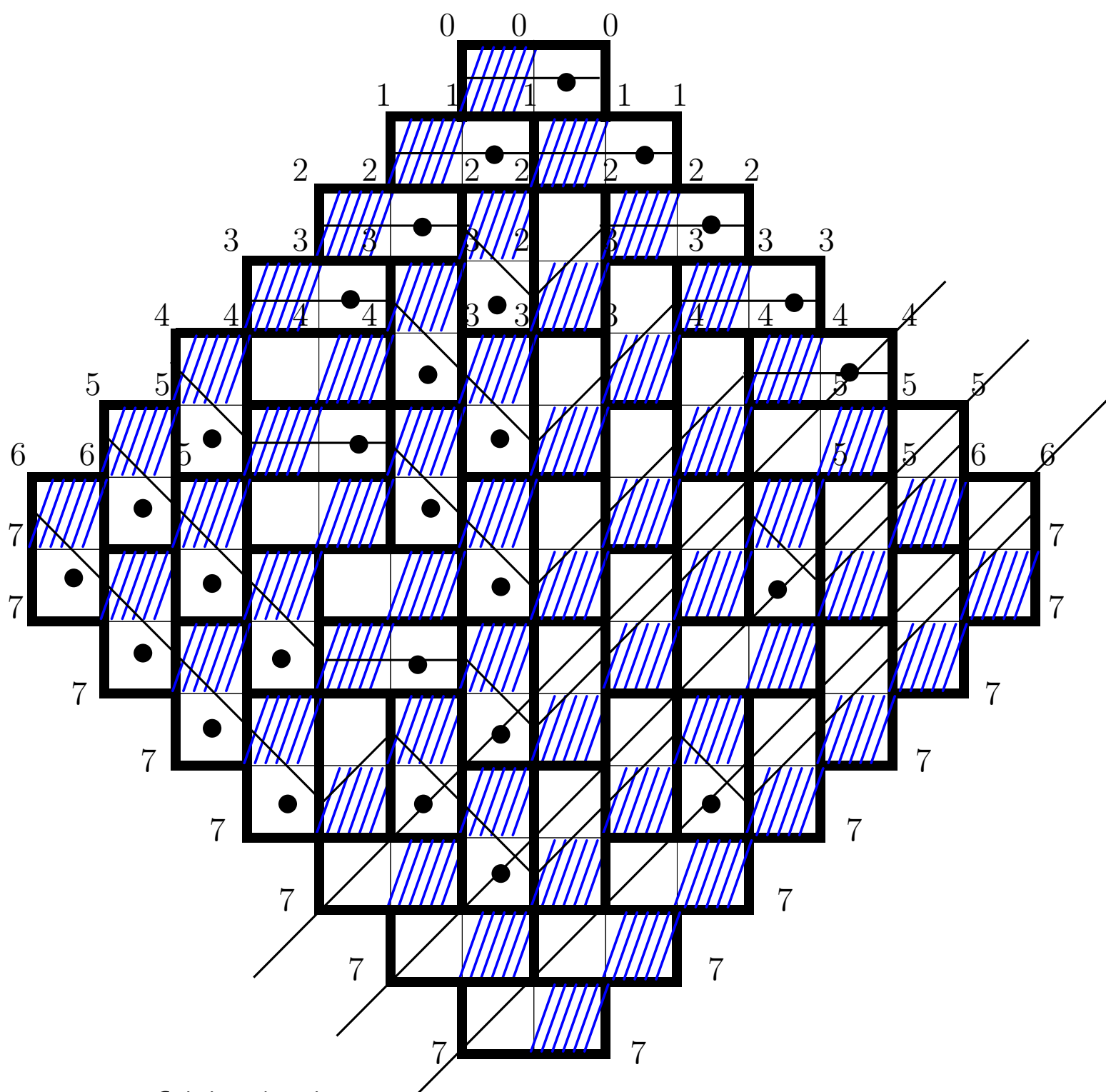


Figure 4. Height function on domino's and level-lines.





Level lines of height $\{\frac{1}{2}, \frac{3}{2}, \dots, n - \frac{1}{2}\}$



Point process of black dots.

A weight on domino's and a probability on domino tilings:

- put the weight $0 < a \leq 1$ on vertical dominoes
- put the weight 1 on horizontal dominoes,

Define:

$$\mathbb{P}(\text{domino tiling } T) = \frac{a^{\#\text{vertical domino's in } T}}{\sum_{\text{all possible tilings } T} a^{\#\text{vertical domino's in } T}}$$

Which probability are we computing?

$\mathbb{P}\{\text{The interval } [k, \ell] \subset \{\xi = 2s\} \text{ contains no dot-particles}\}$

$= \mathbb{P}\{\text{The random surface is flat along the interval } [k, \ell] \subset \{\xi = 2s\}\}$

$= \det \left[\mathbb{1} - \chi_{(k, \ell)}(\eta_1) \mathbb{L}_n^{\text{oneAzT}}(2s, \eta_1; 2s, \eta_2) \chi_{(k, \ell)}(\eta_2) \right]$

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$= \mathbb{P}\{\text{The random surface is flat along the interval } [k, \ell] \subset \{\xi = 2s\}\}$

$= \det \left(\mathbb{1} - \chi_{(k, \ell)}(\eta_1) \mathbb{L}_n^{\text{oneAzt}}(2s, \eta_1; 2s, \eta_2) \chi_{(k, \ell)}(\eta_2) \right)$

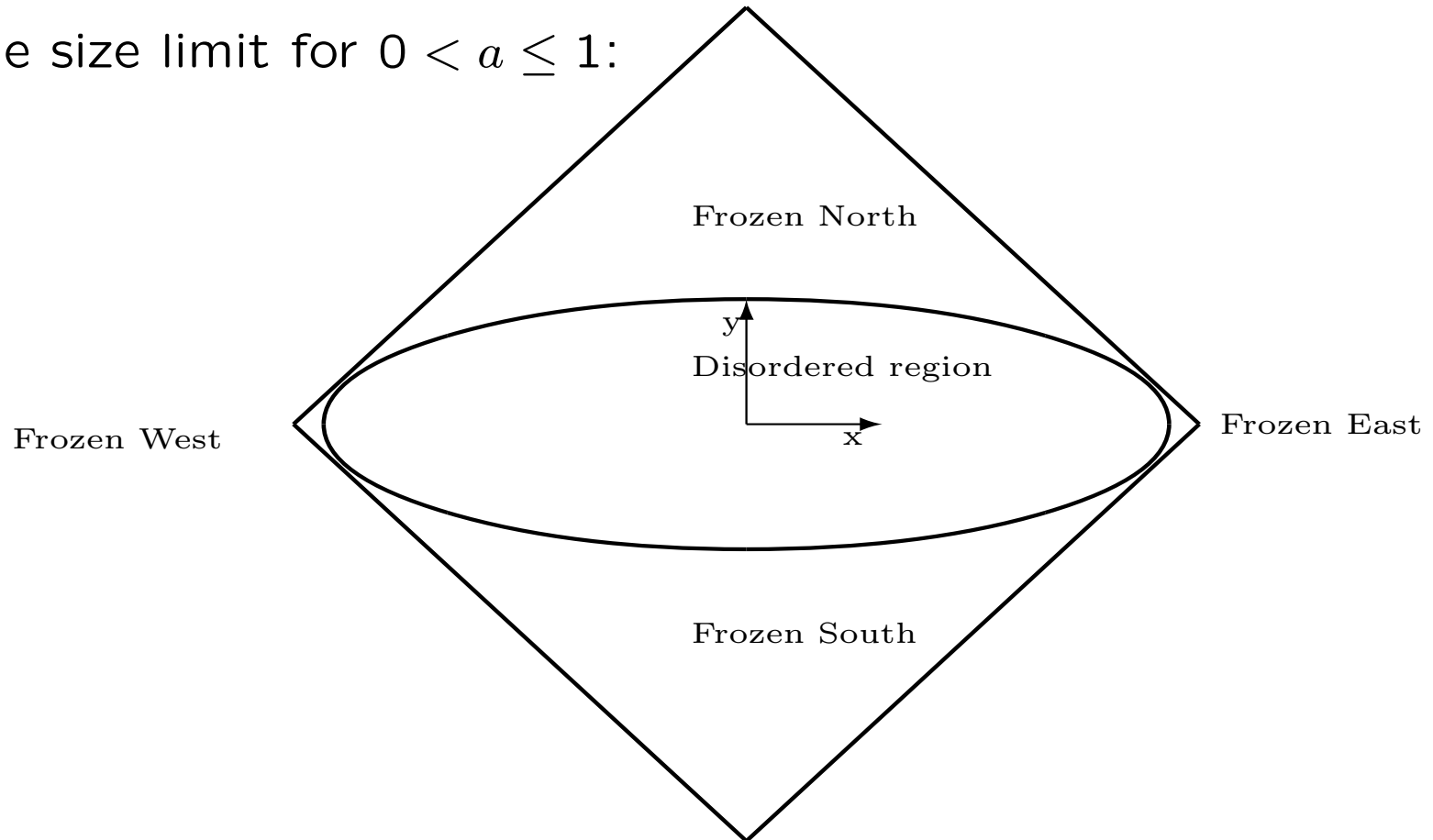
with

$$\begin{aligned} \mathbb{L}_n^{\text{oneAzt}}(\xi_1, \eta_1; \xi_2, \eta_2) = & -\mathbb{1}_{(\xi_1 < \xi_2)} \int_{\Gamma_{0,a}} \frac{dw}{2\pi i} \frac{(1+aw)^{\frac{\eta_1-\eta_2}{2}-1}}{(w-a)^{\frac{\eta_1-\eta_2}{2}+1}} w^{\frac{\xi_1-\xi_2}{2}} \\ & + \int_{\Gamma_{0,a}} \frac{dz}{(2\pi i)^2} \int_{\Gamma_{0,a,z}} \frac{dw}{w-z} \frac{(1+az)^{\frac{\eta_1-1}{2}} (z-a)^{n-\frac{\eta_1+1}{2}}}{(1+aw)^{\frac{\eta_2+1}{2}} (w-a)^{n-\frac{\eta_2-1}{2}}} \frac{w^{n-\frac{\xi_2}{2}}}{z^{n-\frac{\xi_1}{2}}} \end{aligned}$$

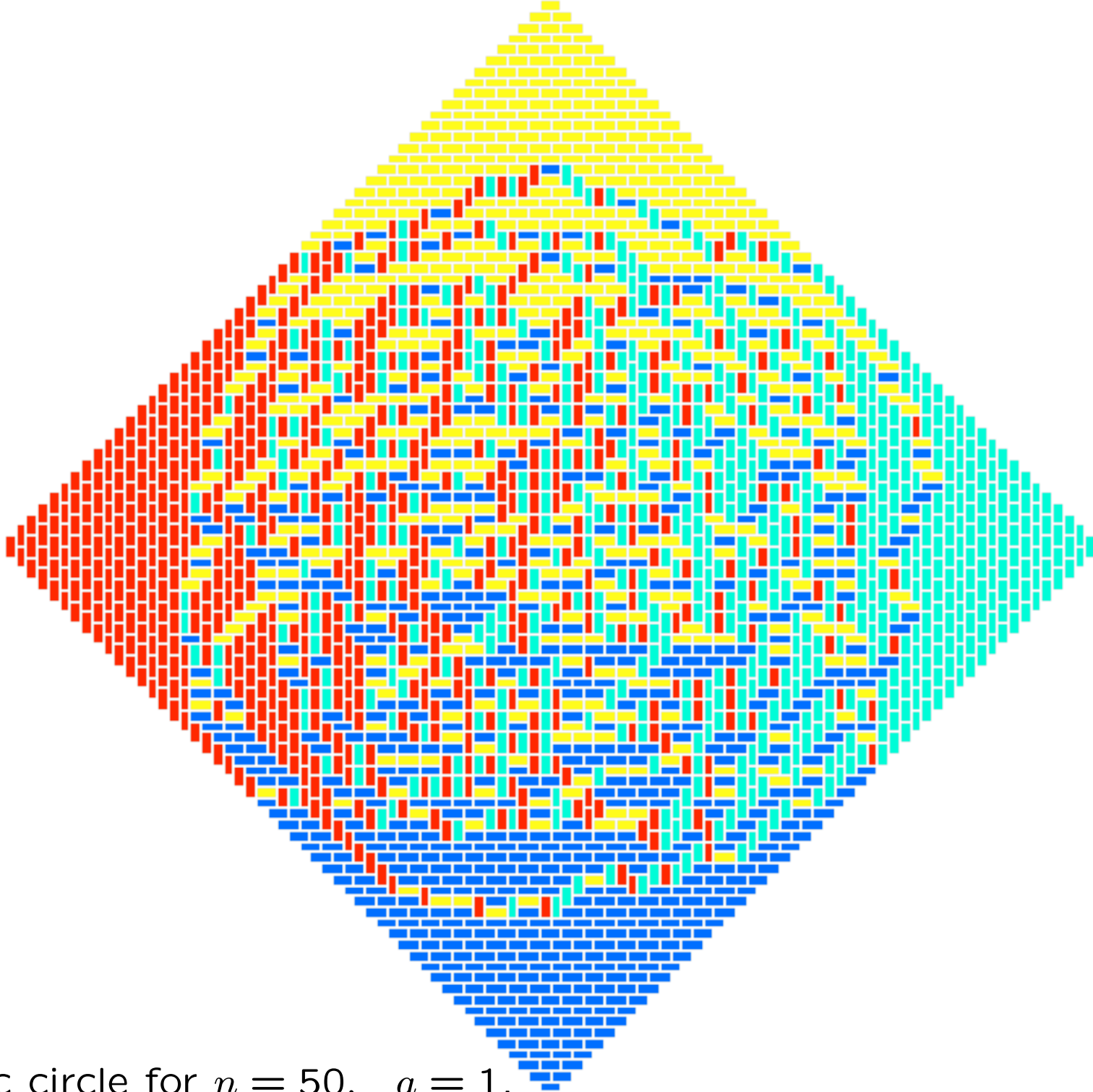
Johansson 03, 05

Letting the size $n \rightarrow \infty$:

In the large size limit for $0 < a \leq 1$:



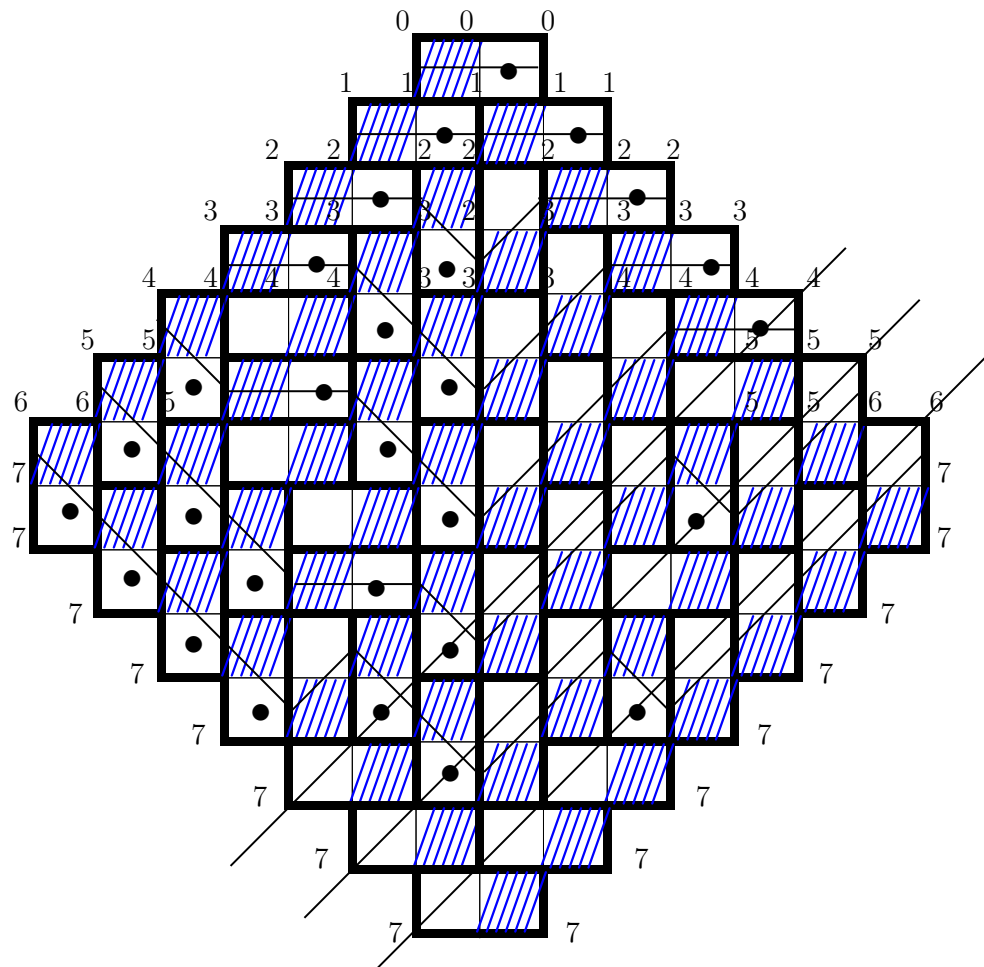
(horizontal domino's are more likely than vertical ones for $0 < a < 1$)



Arctic circle for $n = 50$, $a = 1$.

Determinantal point process for single Aztec diamond:

$$\left\{ \begin{array}{l} \text{Interlacing dots on the} \\ \text{oblique lines for } n \rightarrow \infty \end{array} \right\} \simeq \left\{ \begin{array}{l} \text{good model for the interlacing spectra} \\ \text{of consecutive minors of GUE-matrices} \end{array} \right\}$$



"GUE-minor kernel" ($n \rightarrow \infty$)

(Johansson-Nordenstam '06)

Scaling of the coordinates $\eta_i \in \mathbb{Z}$, $\implies$ coordinates $x, y \in \mathbb{R}$: (Set $a = 1$)

$$\eta_1 = n/2 + x\sqrt{n/2}$$

$$\eta_2 = n/2 + y\sqrt{n/2}$$

LIMITING KERNEL: GUE-minor-Kernel

$$\lim_{n \rightarrow \infty} (\sqrt{n/2}) \mathbb{L}_n^{\text{oneAztec}}(2r, \eta_1; 2s, \eta_2) = \mathbb{K}^{\text{GUE-minor}}(r, x; s, y).$$

with

$$\frac{1}{2} \text{GUE-minor kernel} = \frac{1}{2} \mathbb{K}^{\text{GUE-minor}}(r, x; s, x')$$

$$:= -\mathbb{1}_{r>s} \frac{(2(x-x'))^{r-s-1}}{(r-s-1)!} \mathbb{1}_{x \geq x'} + \int_{\Gamma} \frac{dz}{(2\pi i)^2} \int_L \frac{dw}{w-z} \frac{e^{-z^2+2zx}}{e^{-w^2+2wx'}} \frac{w^s}{z^r}$$

where $\begin{cases} L := 0^+ + i\mathbb{R} \uparrow \text{ to the right of} \\ \Gamma := \text{circle about } 0 \end{cases}$

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Note:

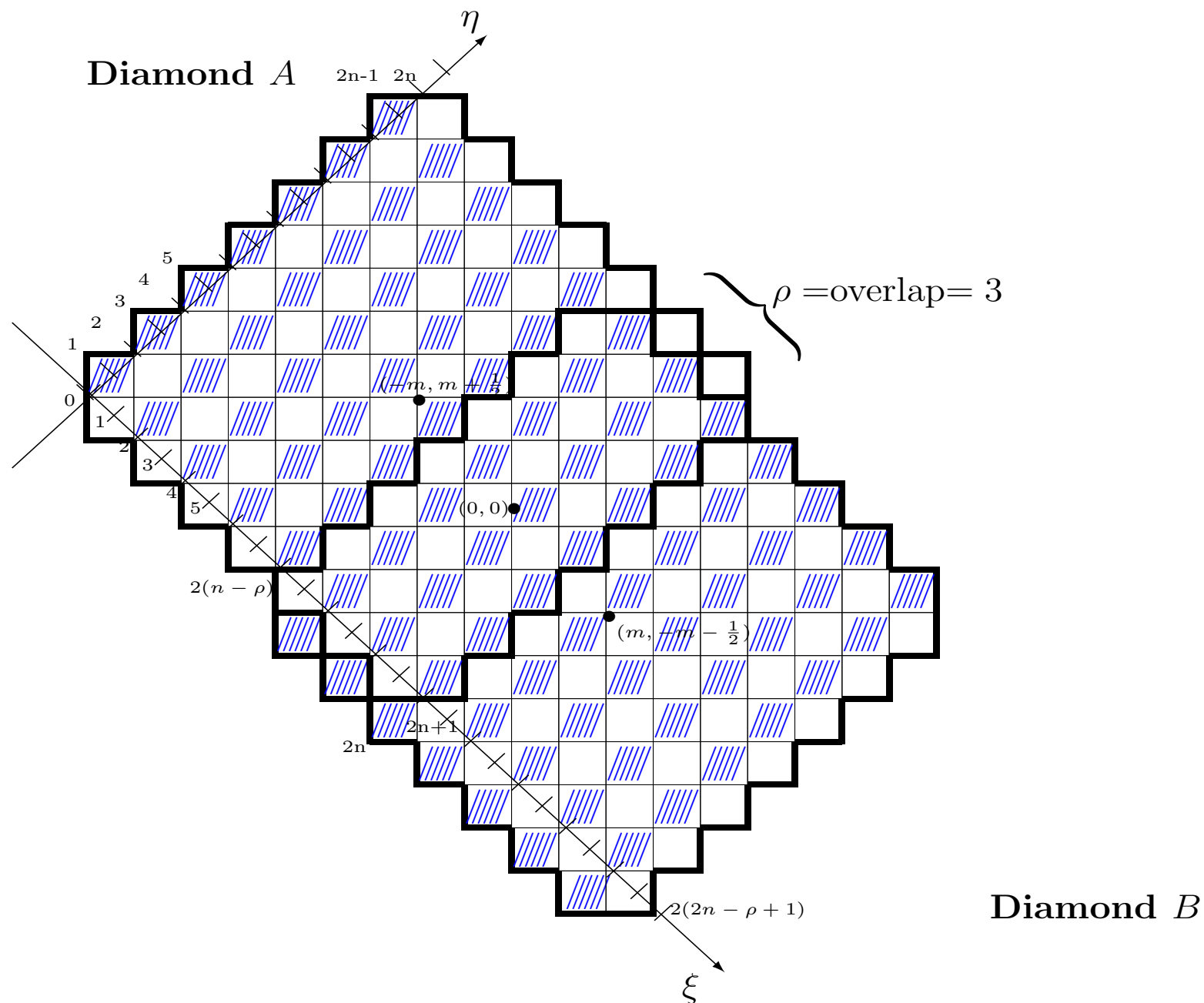
$$\frac{1}{2} \mathbb{K}^{\text{GUE-minor}}(r, x; r, x') = \int_{\Gamma} \frac{dz}{(2\pi i)^2} \int_L \frac{dw}{w-z} \frac{e^{-z^2+2zx}}{e^{-w^2+2wx'}} \frac{w^r}{z^r} = \begin{cases} \text{kernel for} \\ r \times r\text{-GUE} \end{cases}$$

II. Domino tiling of a double Aztec diamond

Adler, Johansson & PvM: *Double Aztec Diamonds and the Tacnode Process*, 2011. (arXiv:1112.5532)

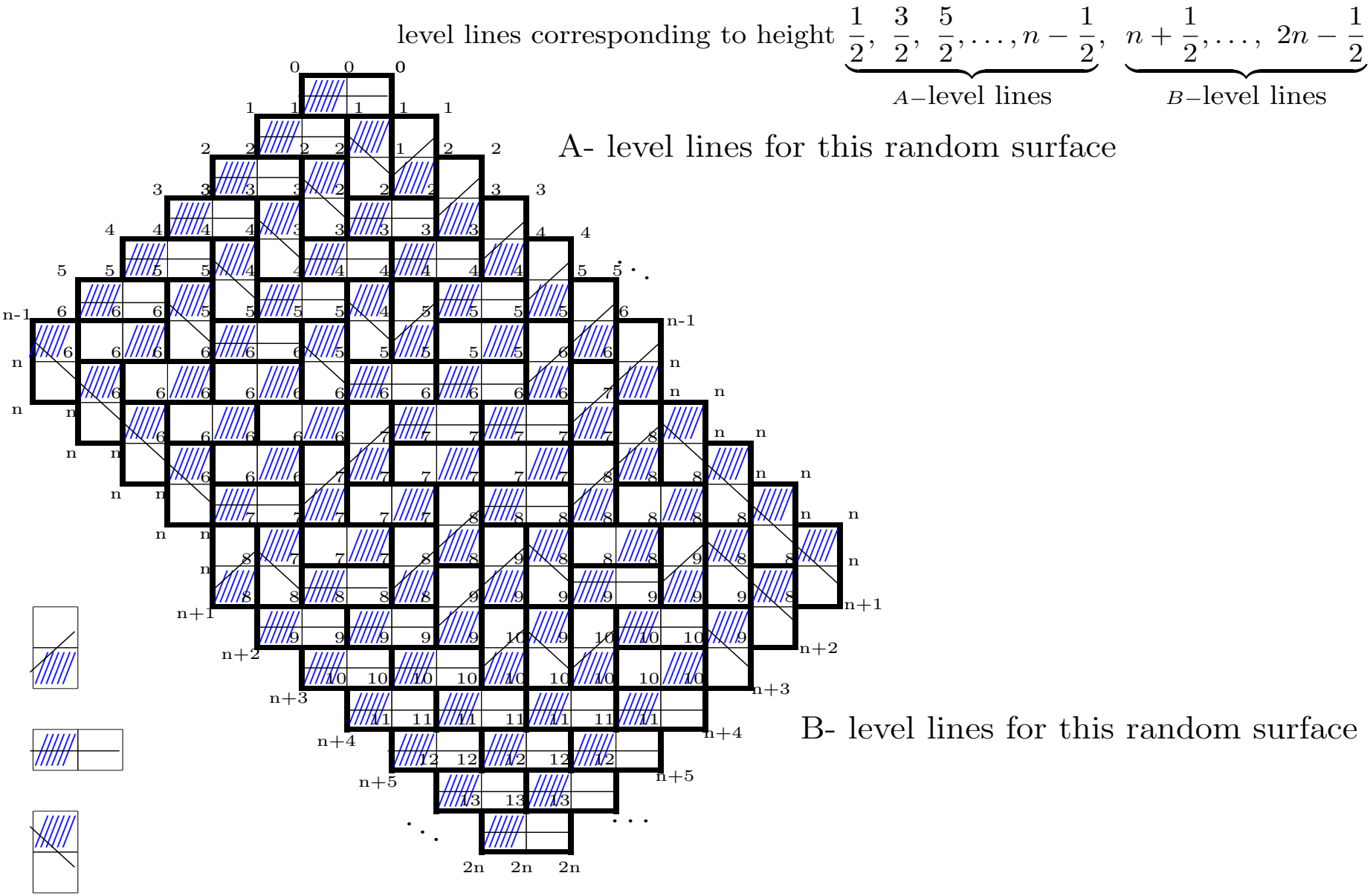
Mark Adler, Sunil Chhita, Kurt Johansson & PvM: *Tacnode GUE-minor Processes and Double Aztec Diamonds*, 2013. (arXiv:1303.5279)

Mark Adler & PvM: *The Double GUE-minor process*, 2013.

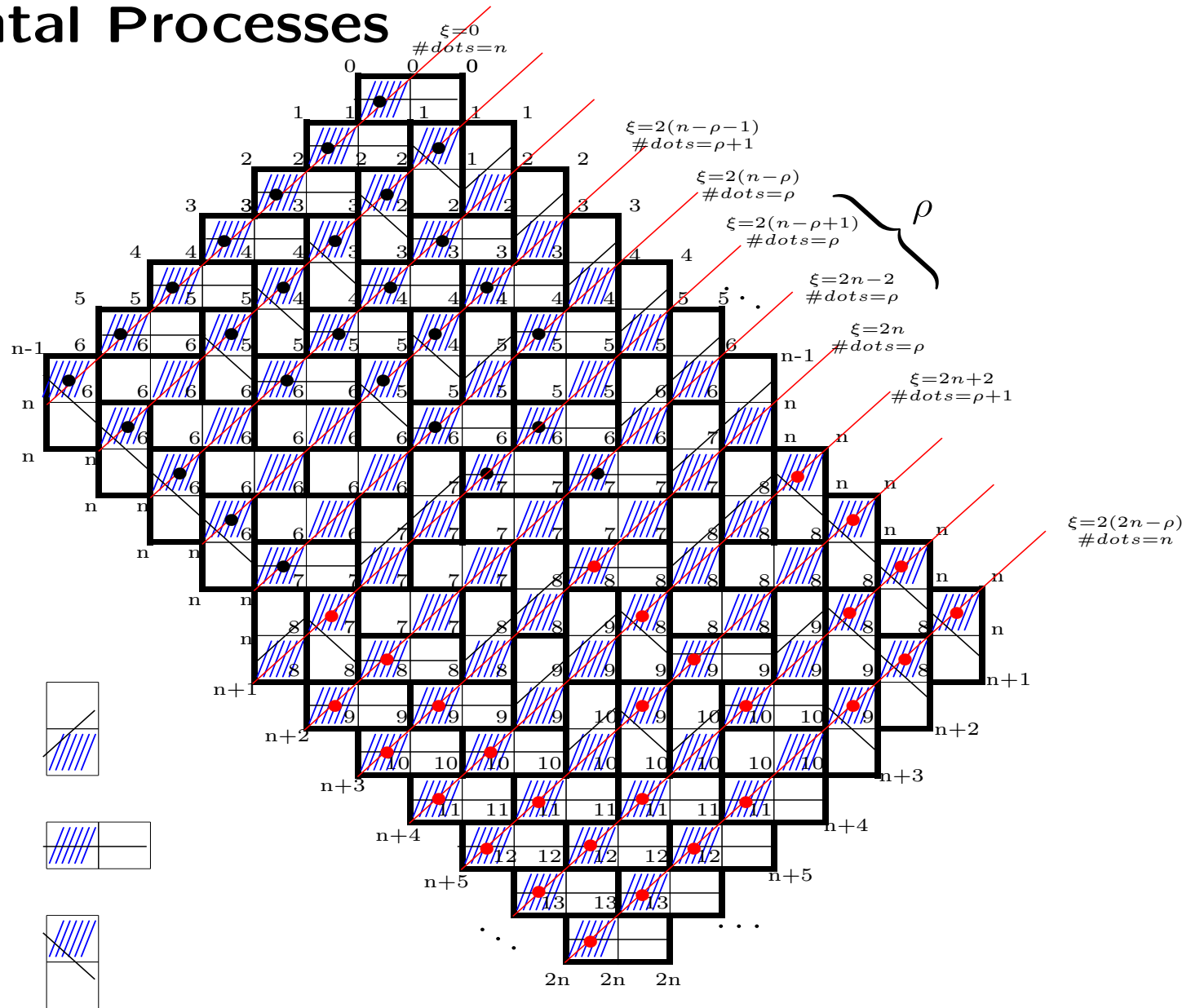


Double Aztec diamond of type $n = 7$, with overlap $= \rho = 3$
 Coordinate system: (ξ, η)

Random cover with domino's: two groups of non-intersecting paths



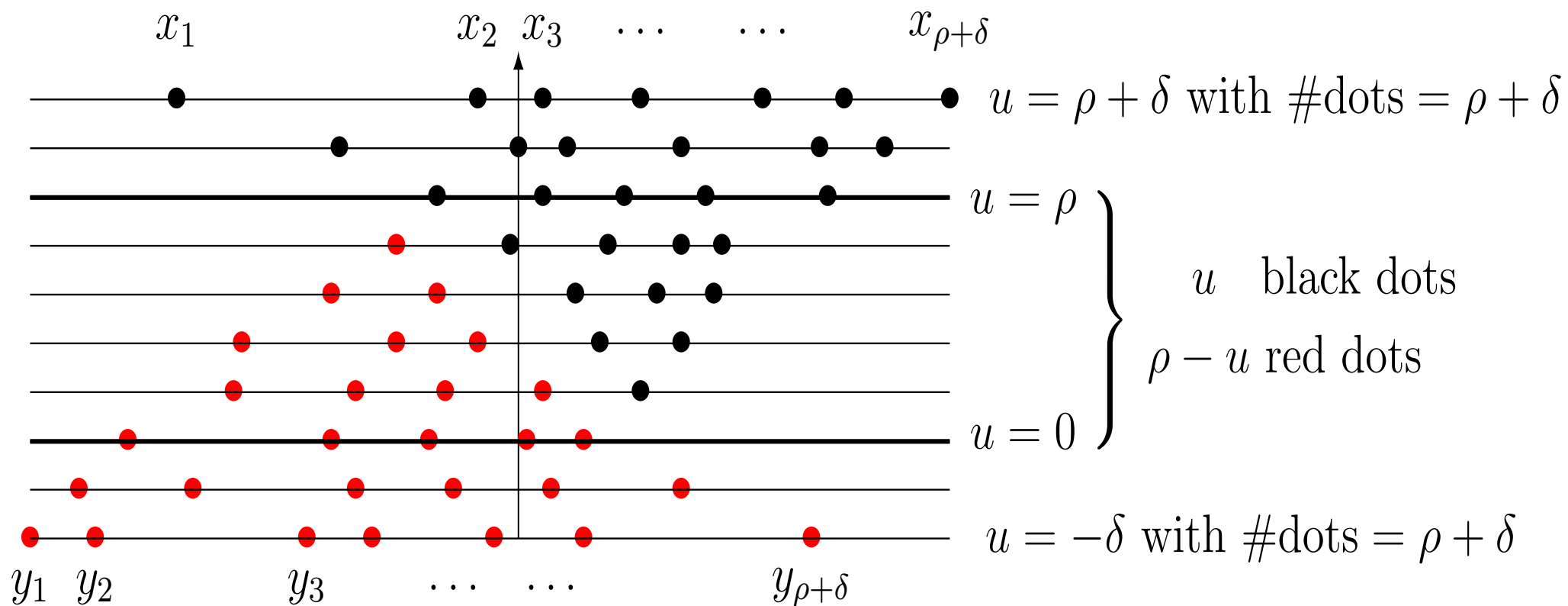
Determinantal Processes



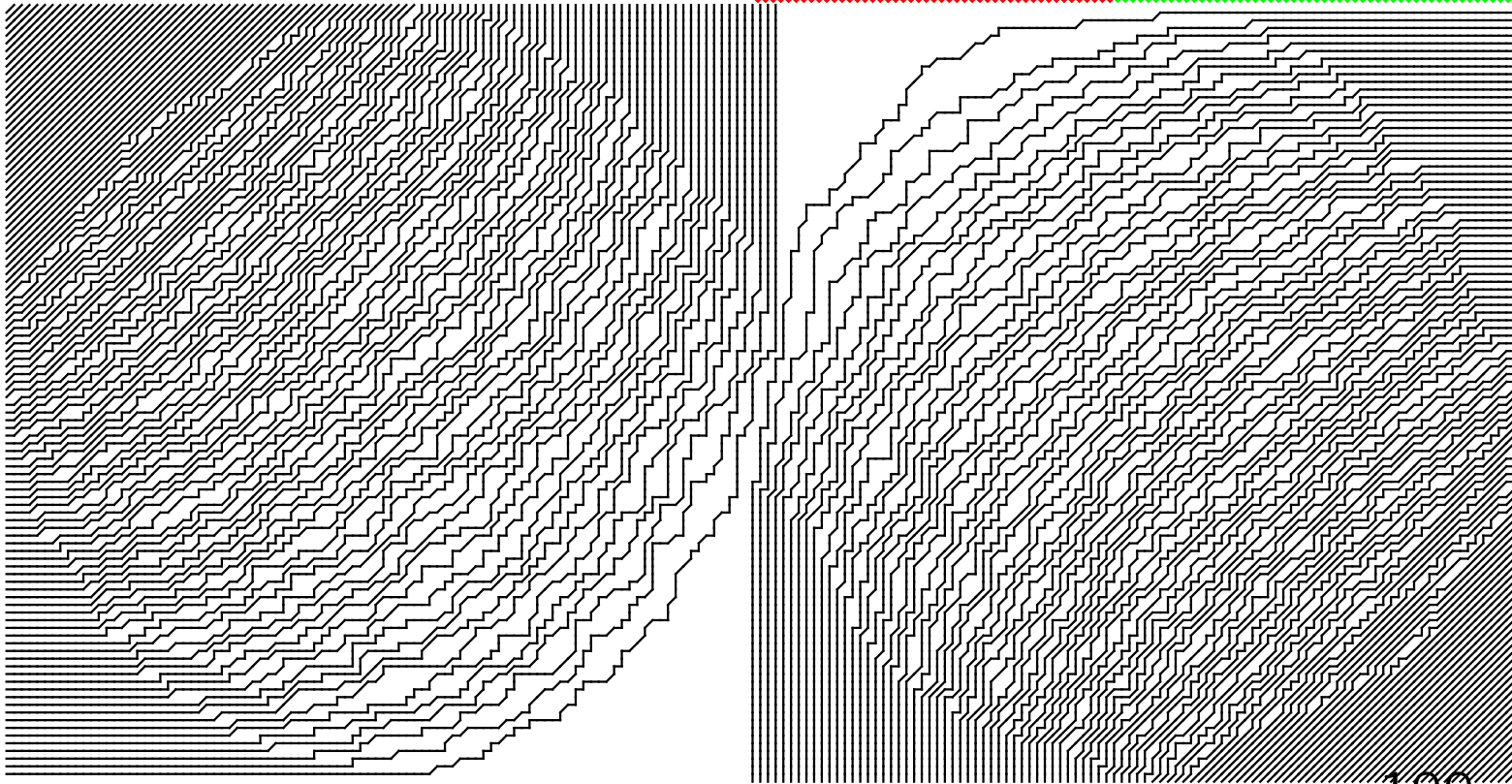
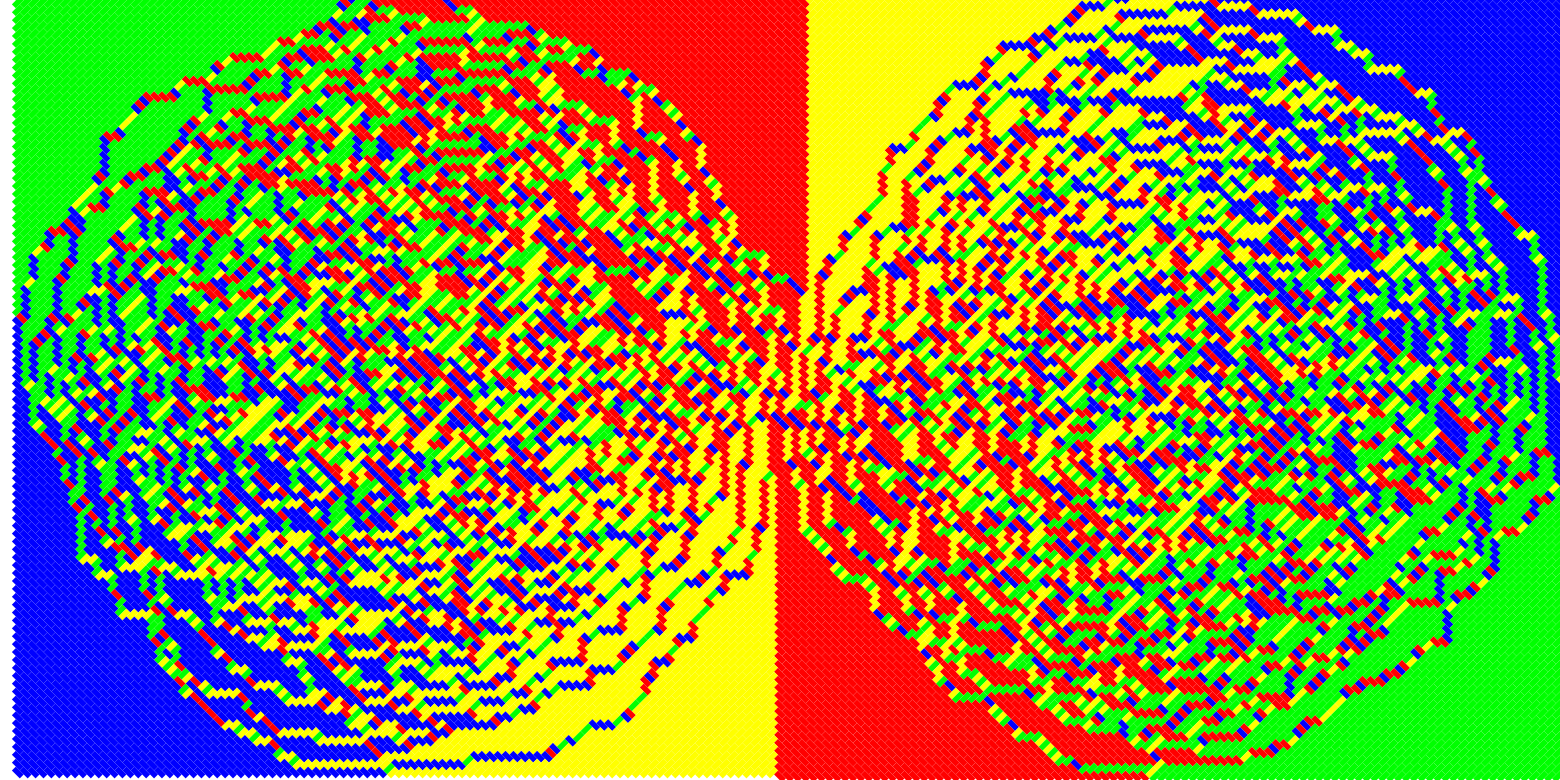
Each time a line $\xi = 2k$ intersects an

- (1) A -level line: put a black dot in that square
 - (2) B -level line: put a red dot in that square
- $\left. \vphantom{\begin{matrix} (1) \\ (2) \end{matrix}} \right\} \textit{interlacing!}$

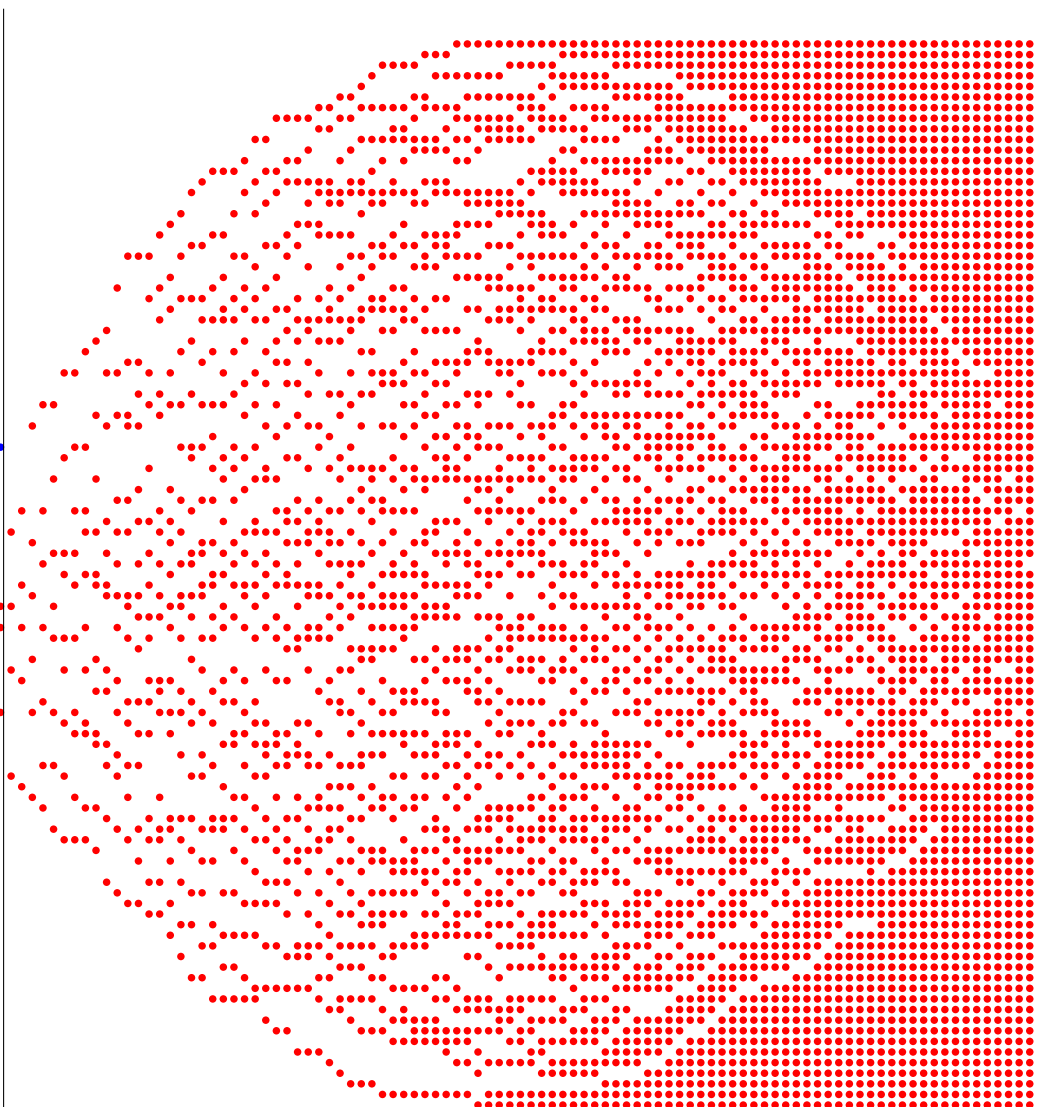
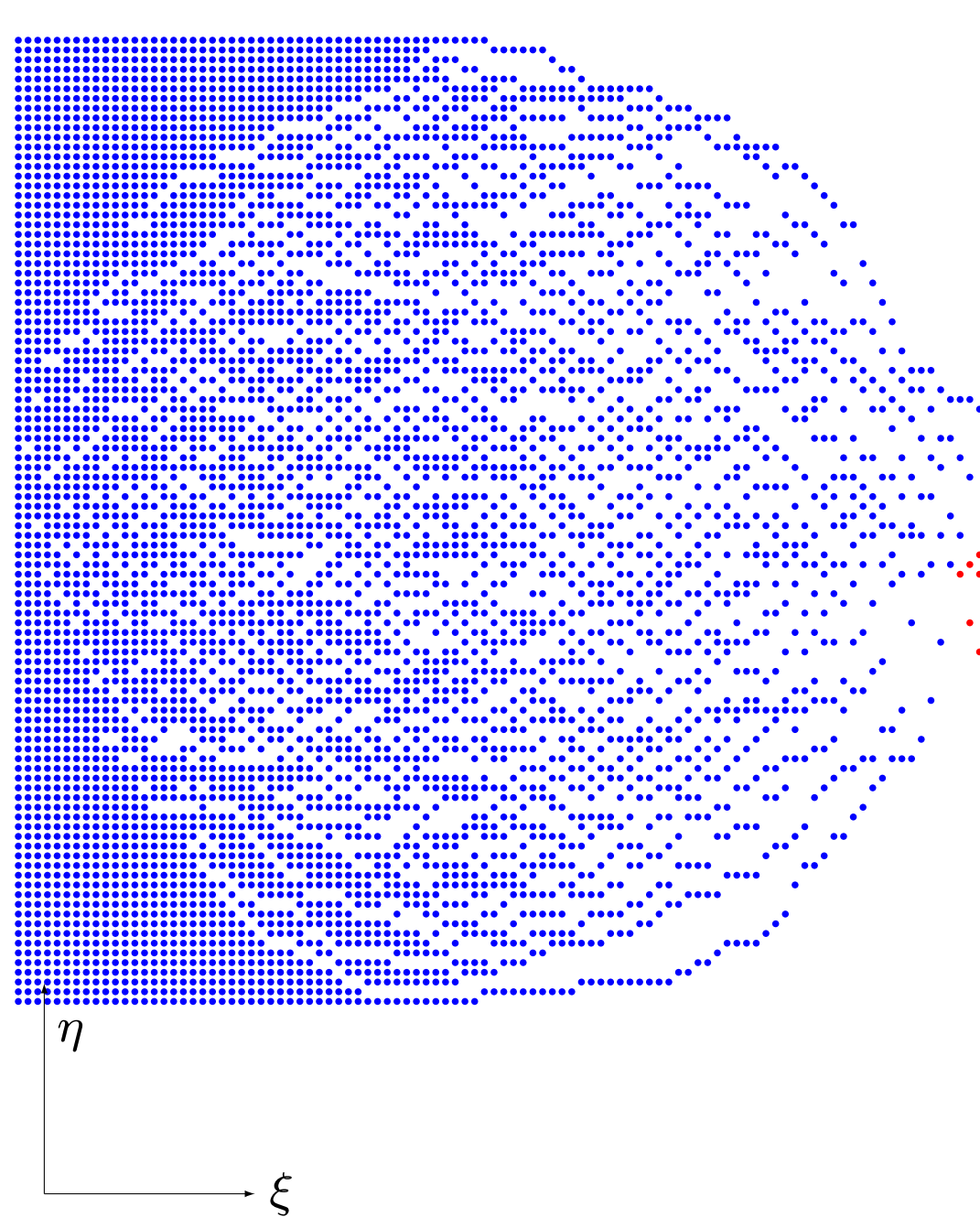
The (discrete) interlacing can schematically be represented by the following picture: (coordinates $2u = 2n - \xi$)



Any such set of (discrete) interlacing black and red dots, respecting the boundary, corresponds to a tiling with domino's !



$n = 100$, $\rho = 4$, $a = 1$



$$n = 100, \rho = 4, a = 1$$

Determinantal point process with kernel $\mathbb{L}_{n,\rho}^{\text{doubleAzt}}(\xi_1, \eta_1; \xi_2, \eta_2)$:

$$\frac{1}{1+a^2} \mathbb{L}_{n,\rho}^{\text{doubleAzt}}(\xi_1, \eta_1; \xi_2, \eta_2) = \mathbb{L}_n^{\text{oneAzt}}(\xi_1, \eta_1; \xi_2, \eta_2)$$

$$- \left\langle ((I - K_0)^{-1} A_{\xi_1, \eta_1})(k), B_{\xi_2, \eta_2}(k) \right\rangle_{\geq n-\rho+1},$$

Via the Kasteleyn matrix !

Determinantal point process with kernel $\mathbb{L}_{n,\rho}^{\text{doubleAzt}}(\xi_1, \eta_1; \xi_2, \eta_2)$:

$$\frac{1}{1+a^2} \mathbb{L}_{n,\rho}^{\text{doubleAzt}}(\xi_1, \eta_1; \xi_2, \eta_2) = \mathbb{L}_n^{\text{oneAzt}}(\xi_1, \eta_1; \xi_2, \eta_2) - \left\langle ((I - K_0)^{-1} A_{\xi_1, \eta_1})(k), B_{\xi_2, \eta_2}(k) \right\rangle_{\geq n-\rho+1},$$

with

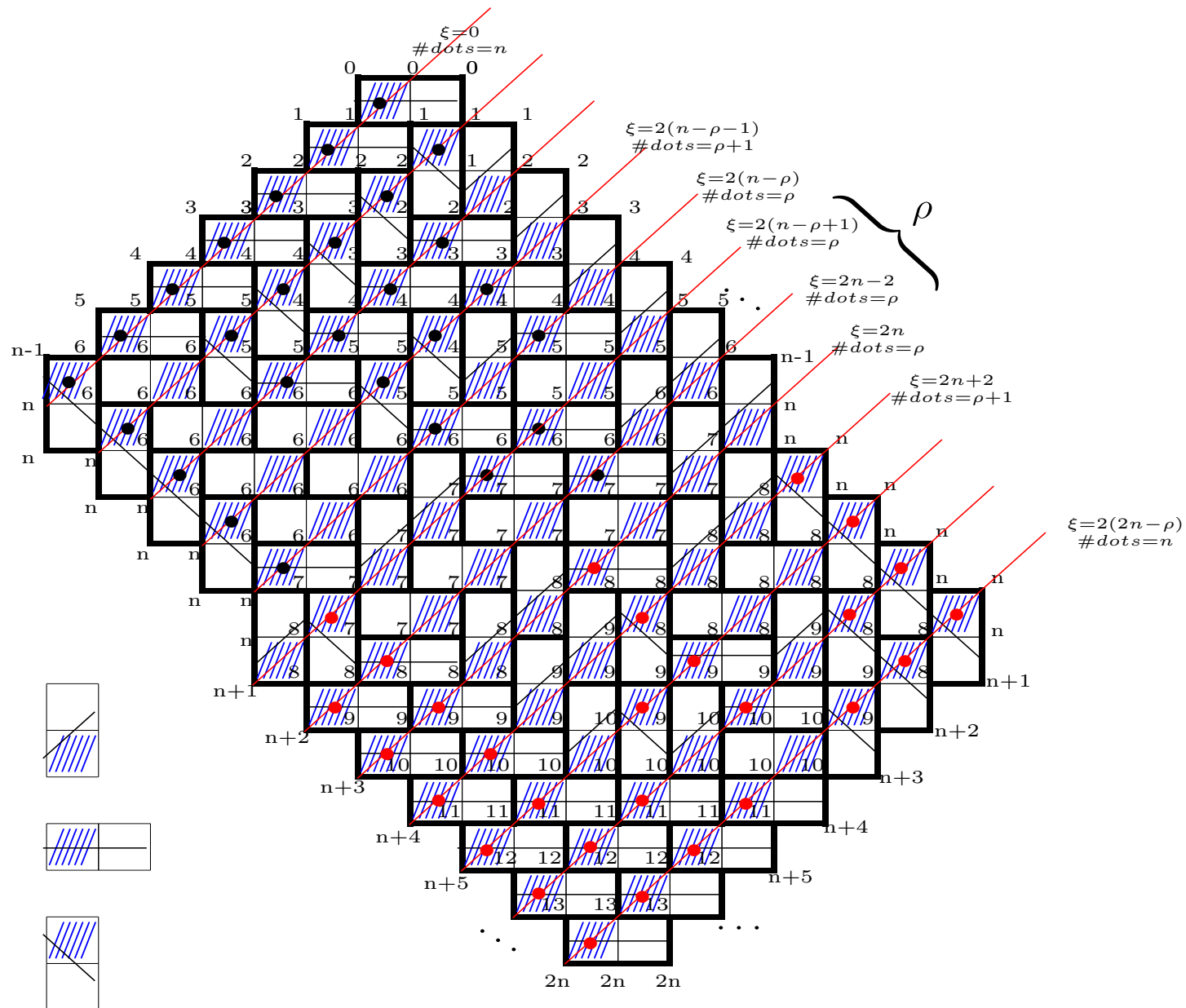
$$K_0(j, k) = \int_{\Gamma_{0,a}} \frac{dw}{(2\pi i)^2} \int_{\Gamma_{0,a,w}} \frac{dz}{z-w} \frac{z^{n-j}}{w^{n+1-k}} \frac{\psi_n(w)}{\psi_n(z)},$$

$$A_{\xi_1, \eta_1}(j) = \int_{\Gamma_{0,a}} \frac{dw}{(2\pi i)^2} \int_{\Gamma_{0,a,w}} \frac{dz}{w-z} \frac{w^{n-j}}{z^{n-\frac{\xi_1}{2}}} \frac{(1+az)^{\frac{\eta_1-1}{2}} (z-a)^{n-\frac{\eta_1+1}{2}}}{\psi_n(w)}$$

$$B_{\xi_2, \eta_2}(k) = \int_{\Gamma_{0,a}} \frac{dz}{(2\pi i)^2} \int_{\Gamma_{0,a,z}} \frac{dw}{w-z} \frac{z^{n-\xi_2/2}}{w^{n+1-k}} \frac{\psi_n(w)}{(1+az)^{\frac{\eta_2+1}{2}} (z-a)^{n-\frac{\eta_2-1}{2}}}$$

$$\psi_n(w) = (1+aw)^n (w-a)^{n+1}$$

III. LETTING THE SIZE $n \rightarrow \infty$,
keeping the overlap finite



Double GUE-minor kernel, as a scaling limit of the double-Aztec kernel

Take scaling limit:

$$\begin{cases} \text{size } n \rightarrow \infty, \\ \rho = \text{fixed overlap of two diamonds}, \\ a = 1 - \beta \sqrt{\frac{2}{n}} = \text{weight of vertical domino's, with } \beta \geq 0. \end{cases}$$

Scaling of the coordinates (ξ, η) , $\implies$ coordinates $(u, y) \in \mathbb{Z} \times \mathbb{R}$:

$$\xi_i = 2n - 2u_i, \quad \eta = n + [y\sqrt{2n}] - 1.$$

LIMITING KERNEL: Double GUE-Kernel

$$\begin{aligned} \lim_{n \rightarrow \infty} (-a)^{\frac{1}{2}(\eta_1 - \eta_2)} (-\sqrt{n/2})^{\frac{1}{2}(\xi_1 - \xi_2)} \mathbb{L}_{n,\rho}^{\text{doubleAzt}}(\xi_1, \eta_1; \xi_2, \eta_2) \frac{\Delta \eta_2}{2} \\ = \mathbb{K}_{\beta,\rho}^{\text{dGUE}}(u_1, y_1; u_2, y_2) dy_2. \end{aligned}$$

Double GUE-minor kernel, two parameters: $\begin{cases} \beta = \text{speed at which } a \rightarrow 1 \\ \rho = \text{overlap of two diamonds} \end{cases}$

$$\begin{aligned} \frac{1}{2}\mathbb{K}_{\beta,\rho}^{\text{dGUE}}(u_1,y_1;u_2,y_2) &= \frac{1}{2}\mathbb{K}^{\text{GUEminor}}(u_1,\beta-y_1; \, u_2,\beta-y_2) \\ &\quad + \left\langle (\mathbb{1}-\mathcal{K}^\beta(\lambda,\kappa))^{-1} \, \mathcal{A}_{u_1}^{\beta,y_1-\beta}(\kappa), \mathcal{B}_{u_2}^{\beta,y_2-\beta}(\lambda) \right\rangle_{\geq -\rho} \end{aligned}$$

with $u_i \in \mathbb{Z}, \quad y_i \in \mathbb{R}, \quad \lambda, \mu \in \mathbb{Z}$

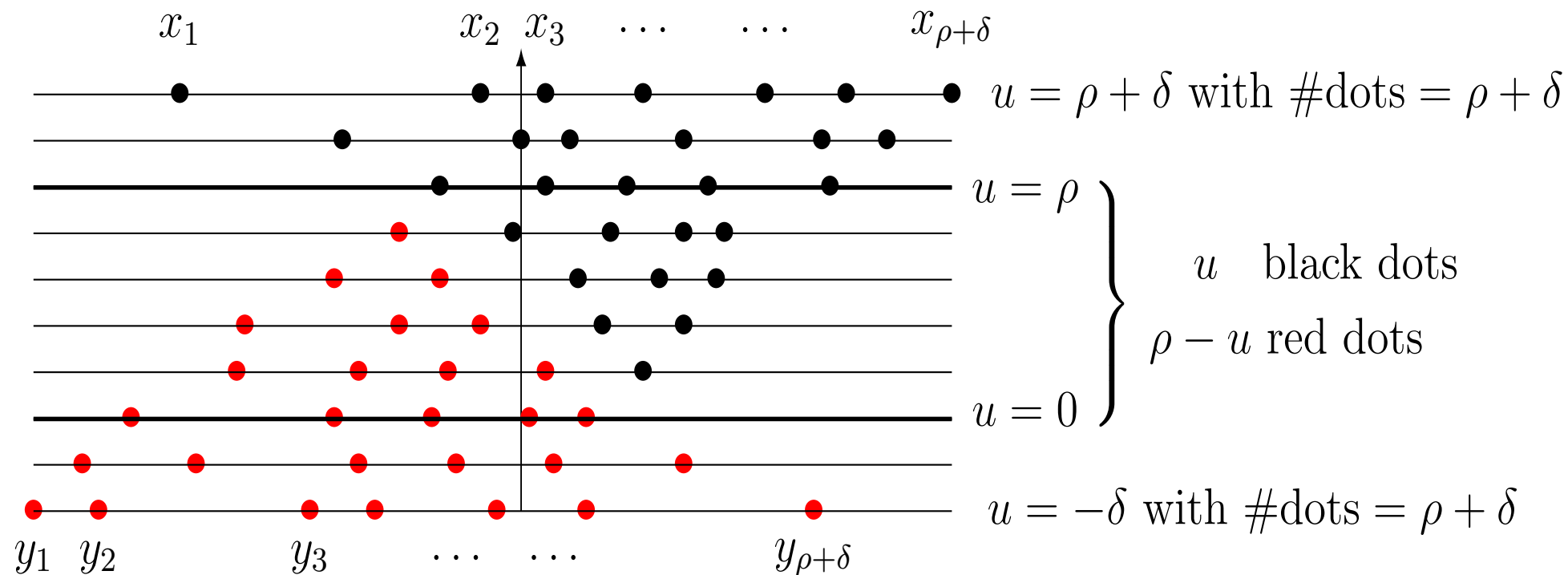
Double GUE-minor kernel, two parameters: $\begin{cases} \beta = \text{speed at which } a \rightarrow 1 \\ \rho = \text{overlap of two diamonds} \end{cases}$

$$\begin{aligned} \frac{1}{2} \mathbb{K}_{\beta, \rho}^{\text{dGUE}}(u_1, y_1; u_2, y_2) &= \frac{1}{2} \mathbb{K}^{\text{GUEminor}}(u_1, \beta - y_1; u_2, \beta - y_2) \\ &\quad + \left\langle (\mathbb{1} - \mathcal{K}^\beta)^{-1} \mathcal{A}_{u_1}^{\beta, y_1 - \beta}(\lambda), \mathcal{B}_{u_2}^{\beta, y_2 - \beta}(\lambda) \right\rangle_{\geq -\rho} \\ &\quad \uparrow \end{aligned}$$

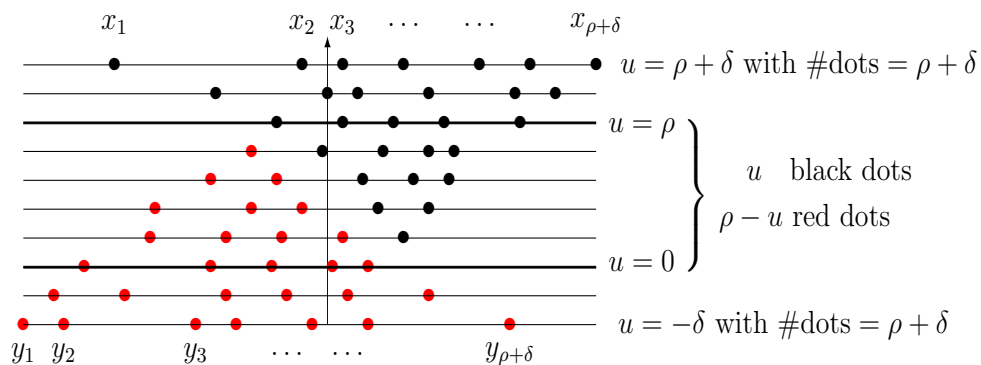
(Finite rank perturbation of GUE-minor kernel!)

$$\begin{aligned} \mathcal{K}^{(\beta)}(\lambda, \kappa) &:= \oint_{\Gamma} \frac{d\zeta}{(2\pi i)^2} \int_L \frac{d\omega}{\omega - \zeta} \frac{e^{-2\zeta^2 + 4\beta\zeta}}{e^{-2\omega^2 + 4\beta\omega}} \frac{\zeta^\kappa}{\omega^{\lambda+1}}, \quad \lambda, \kappa \in \mathbb{Z} \\ \mathcal{A}_v^{\beta, y}(\kappa) &:= \oint_{\Gamma} \frac{d\zeta}{(2\pi i)^2} \int_L \frac{d\omega}{\zeta - \omega} \frac{e^{-\zeta^2 - 2y\zeta}}{e^{-2\omega^2 + 4\beta\omega}} \frac{\zeta^{-v}}{\omega^{\kappa+1}} + \int_L \frac{d\zeta}{2\pi i} \frac{e^{\zeta^2 - 2\zeta(y+2\beta)}}{\zeta^{v+\kappa+1}} \\ \mathcal{B}_u^{\beta, y}(\lambda) &:= \oint_{\Gamma} \frac{d\zeta}{(2\pi i)^2} \int_L \frac{d\omega}{\zeta - \omega} \frac{e^{-2\zeta^2 + 4\zeta\beta}}{e^{-\omega^2 - 2\omega y}} \frac{\zeta^\lambda}{\omega^{-u}} + \oint_{\Gamma} \frac{d\omega}{2\pi i} \frac{\omega^{u+\lambda}}{e^{\omega^2 - 2\omega(y+2\beta)}}. \end{aligned}$$

$$\text{where } \begin{cases} L := 0^+ + i\mathbb{R} \uparrow \\ \Gamma := \text{circle about } 0 \end{cases}$$



In the limit, when size $\rightarrow \infty$: *Continuous interlacing* of red and blue dots for $\rho = 5$ and $\delta = 2$.



Level $u = -\delta$ at a distance $-\delta$ below the overlap:

density of particles is given by: $(\mathbf{y} = (y_1, \dots, y_{\rho+\delta}))$

$$\det \left(\mathbb{K}_{\beta, \rho}^{\text{dGUE}}(-\delta, y_i; -\delta, y_j) \right)_{1 \leq i, j \leq \rho+\delta} = C_{\rho, \delta} \Delta_{\rho+\delta}(\mathbf{y}) \widetilde{\Delta}_{\rho+\delta}(\mathbf{y}) \prod_1^{\rho+\delta} \frac{e^{-(y_i + \beta)^2}}{\sqrt{\pi}}$$

where

$$\Delta_{\rho+\delta}(\mathbf{y}) = \det \begin{pmatrix} 1 \\ y \\ \vdots \\ y^{\delta-1} \\ y^{\delta} \\ \vdots \\ y^{\delta+\rho-1} \end{pmatrix}_{y=y_1, \dots, y_{\rho+\delta}} \quad \widetilde{\Delta}_{\rho+\delta}(\mathbf{y}) := \det \begin{pmatrix} 1 \\ y \\ \vdots \\ y^{\delta-1} \\ \Phi_{\delta}(\beta - y) \\ \vdots \\ \Phi_{\delta+\rho-1}(\beta - y) \end{pmatrix}_{y=y_1, \dots, y_{\rho+\delta}}$$

Define for $n \geq 0$

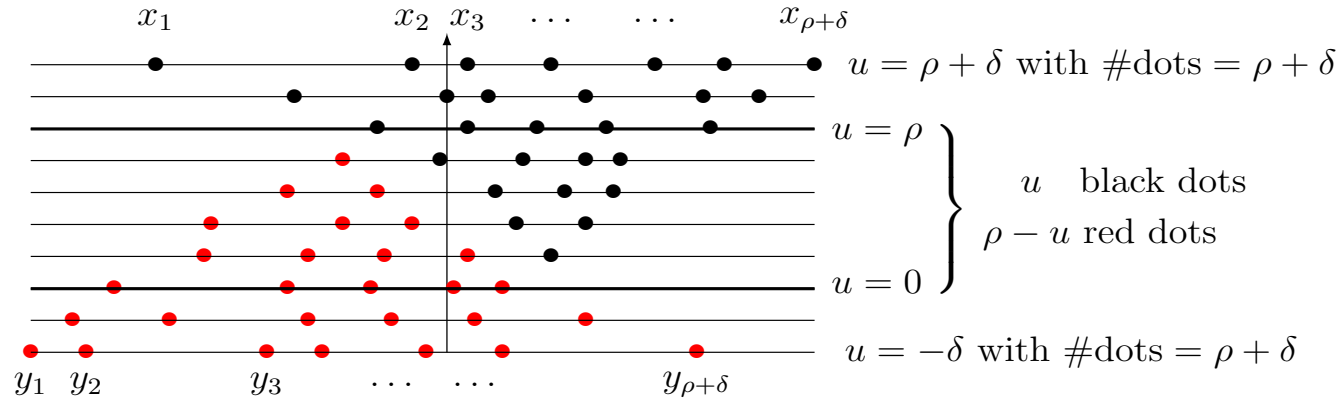
$$\Phi_n(y) := P_n(y) \left(\frac{1}{\sqrt{\pi}} \int_{-\infty}^y e^{-\xi^2} d\xi \right) + Q_{n-1}(y) \left(\frac{1}{\sqrt{\pi}} e^{-y^2} \right)$$

with

$$P_n(y) := \frac{1}{n! \, i^n} H_n(iy),$$

$Q_n(y) :=$ polynomials, defined by recursion relation

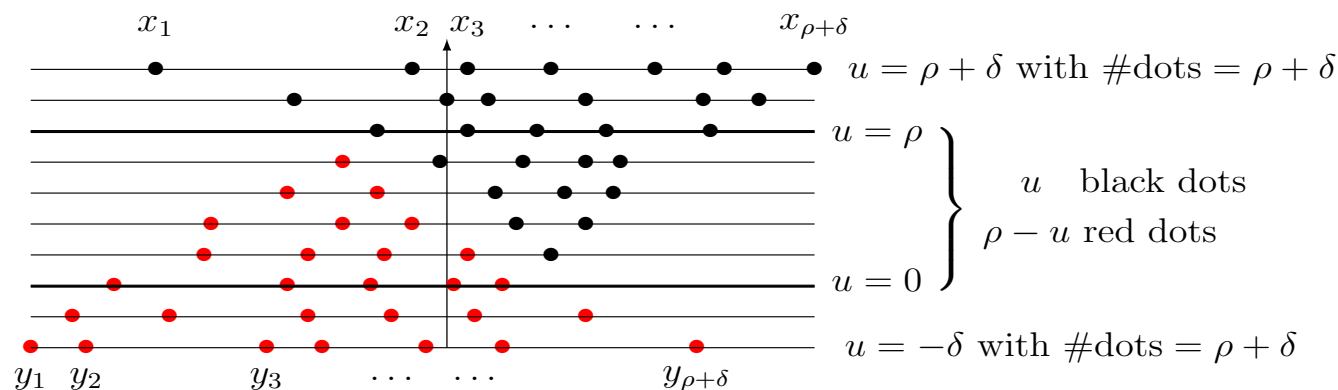
$$\frac{1}{2}(n+1)Q_n(y) = yQ_{n-1}(y) + Q_{n-2}(y) \text{ with } Q_0 = 1 \text{ and } Q_{-1} = 0$$



At levels $-\delta$ and $\rho + \delta$ (below and above the overlap) :
joint density of particles at levels is given by:

$$\mathbb{P} \left(\bigcap_1^{\rho+\delta} \{x_i \in dx_i, y_i \in dy_i\} \right) := dx dy$$

$$\times \det \begin{pmatrix} \left(\frac{1}{2} \mathbb{K}_{\beta, \rho}^{\text{dGUE}}(-\delta, y_i; -\delta, y_j) \right)_{1 \leq i, j \leq \rho+\delta} & \left(\frac{1}{2} \mathbb{K}_{\beta, \rho}^{\text{dGUE}}(-\delta, y_i; \rho + \delta, x_j) \right)_{1 \leq i, j \leq \rho+\delta} \\ \left(\frac{1}{2} \mathbb{K}_{\beta, \rho}^{\text{dGUE}}(\rho + \delta, x_i; -\delta, y_j) \right)_{1 \leq i, j \leq \rho+\delta} & \left(\frac{1}{2} \mathbb{K}_{\beta, \rho}^{\text{dGUE}}(\rho + \delta, x_i; \rho + \delta, x_j) \right)_{1 \leq i, j \leq \rho+\delta} \end{pmatrix}$$



- Two level density at levels $-\delta$ and $\rho + \delta$

$$\mathbb{P} \left(\bigcap_1^{\rho+\delta} \{x_i \in dx_i, y_i \in dy_i\} \right) = C \left(\Delta^2(\mathbf{x}) \prod_{i=1}^{\rho+\delta} e^{-(x_i - \beta)^2} \right) \left(\Delta^2(\mathbf{y}) \prod_{i=1}^{\rho+\delta} e^{-(y_i + \beta)^2} \right) \\ \times \frac{\text{Vol}(\mathcal{C}(\mathbf{x}, \mathbf{y}))}{\text{Vol}(\mathcal{C}(\mathbf{x})) \text{Vol}(\mathcal{C}(\mathbf{y}))} \prod_1^{\rho+\delta} dx_i dy_i$$

where $\text{Vol}(\mathcal{C}(\mathbf{x})) = \text{Volume}(\text{cone about } \mathbf{x} = (x_1, \dots, x_{\rho+\delta}))$

$$\text{Vol}(\mathcal{C}(\mathbf{x}, \mathbf{y})) = \text{Volume} \left(\begin{array}{c} \text{"double cone"} \\ \text{about } (\mathbf{x}, \mathbf{y}) \end{array} \right) \\ = \int_{\mathbb{R}^{(\rho+\delta)(\rho+\delta-1)}} \mathbb{1}_{\mathbf{x} \prec \mathbf{z}^{(-\delta+1)} \prec \dots \prec \mathbf{z}^{\rho+\delta-1} \prec \mathbf{y}} (\mathbf{z}) d\mathbf{z}$$

For real physics application of the generic points of the boundary between the frozen and the stochastic region, see:

K. A. Takeuchi & M. Sano: *Evidence for geometry-dependent universal fluctuations of the Kardar-Parisi-Zhang interfaces in liquid-crystal turbulence*, J. Stat. Phys. 147, 853-890 (2012).

<http://dx.doi.org/10.1007/s10955-012-0503-0>